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# **Multiple Phase Screen (MPS) Calculation of Two-way Spherical Wave Propagation in the Ionosphere**

## **Ionospheric Effects Symposium**

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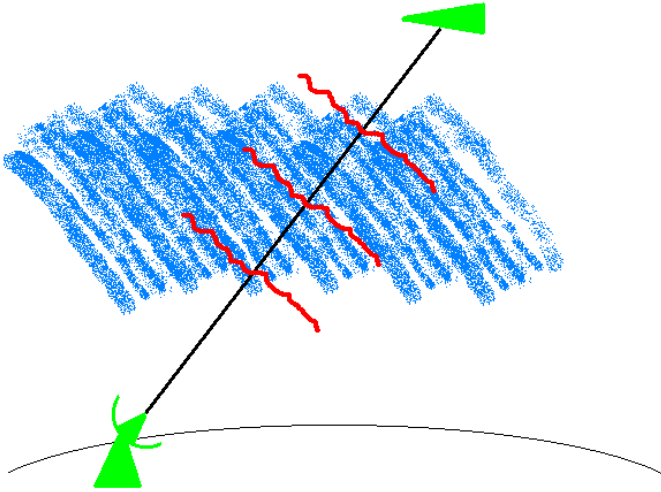
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- **Introduction**
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# Introduction (1/2)

## MPS Signal Generation



- **Parabolic Wave Equation for E-field**

$$\left(\frac{1}{z'^2}\right) \left(\frac{\partial^2 U}{\partial \theta^2} + \frac{\partial^2 U}{\partial \phi^2}\right) - i2k \frac{\partial U}{\partial z'} + k^2 \beta \epsilon U = 0$$

↑  
Diffraction  
term

↑  
Source  
term

### Solution Method:

- **Collapse ionospheric structure to multiple thin phase-changing screens with free space between**
- **At phase screen, neglect diffraction term**
- **Between screens, the PWE is source free, so can solve by Fourier Transform method**
- **Solution U is the single-frequency transfer function. U is the Fourier transform of the impulse response function.**



- **Impulse response function**
  - **Convolve the impulse response function with the transmitted waveform to obtain the received, disturbed waveform**
- **Two methods to calculate the impulse response function:**
  - **Statistical techniques:**
    - > **Techniques based on the mutual coherence function (MCF)**
    - > **Starting point is the analytic solution for the two-frequency, two-time, two-position MCF (the correlation function of the propagating electric field)**
    - > **Theoretical calculation requires strong scattering,  $S_4$  equal to unity, phase structure function must be quadratic, signal bandwidth is small, structure is homogeneous.**
    - > **Limitations never fully studied**
    - > **Previously the choice for most receiver testing because of speed and relative simplicity. But, still in use now for strategic systems**
  - **Multiple phase screen (MPS) techniques**
    - > **Most accurate technique available. Starting point is a realization of the in-situ electron density. None of the limitations above apply.**



## Scalar Helmholtz equation

$$\nabla^2 \psi + k^2 (1 + \beta \epsilon) \psi = 0$$

where

$$\beta = \frac{-\omega_p^2}{\omega^2 - \omega_p^2}$$

$$\epsilon = \frac{\Delta N_e}{\langle N_e \rangle}$$

$$\omega_p^2 = 4\pi c^2 r_e \langle N_e \rangle$$

$$k = \frac{2\pi}{\lambda}$$



**Substitute the parabolic approximation for a spherical wave**

$$\psi(x, y, z) = \frac{U(x, y, z)}{(z - z_t)} \times \exp \left\{ -ik \left( z - z_t + \frac{(x - x_t)^2 + (y - y_t)^2}{2(z - z_t)} \right) \right\}$$

**Make the substitutions**

$$\theta = \frac{(x - x_t)}{z'}; \quad \phi = \frac{(y - y_t)}{z'}; \quad z' = z - z_t$$

**To obtain the final parabolic wave equation (PWE)**

$$\left( \frac{1}{z'^2} \right) \left( \frac{\partial^2 U}{\partial \theta^2} + \frac{\partial^2 U}{\partial \phi^2} \right) - i2k \frac{\partial U}{\partial z'} + k^2 \beta \epsilon U = 0$$

**Propagation through a phase screen: solve PWE with diffraction term set to zero**

$$U \left( \theta, \frac{\Delta z'}{2} \right) = U \left( \theta, -\frac{\Delta z'}{2} \right) \exp \left\{ -ik \int_{-\frac{\Delta z'}{2}}^{\frac{\Delta z'}{2}} \Delta n(\theta, \zeta) d\zeta \right\}$$



**Free-space propagation between phase screens: set source term to zero and solve remaining equation via FFTs**

$$U(\theta, z') = \int_{-\infty}^{\infty} \hat{U}(q_\theta, z') \exp(i2\pi q_\theta \theta) dq_\theta$$

$$\hat{U}(q_\theta, z') = \int_{-\infty}^{\infty} U(\theta, z') \exp(-i2\pi q_\theta \theta) d\theta$$

**The solution for free-space propagation is**

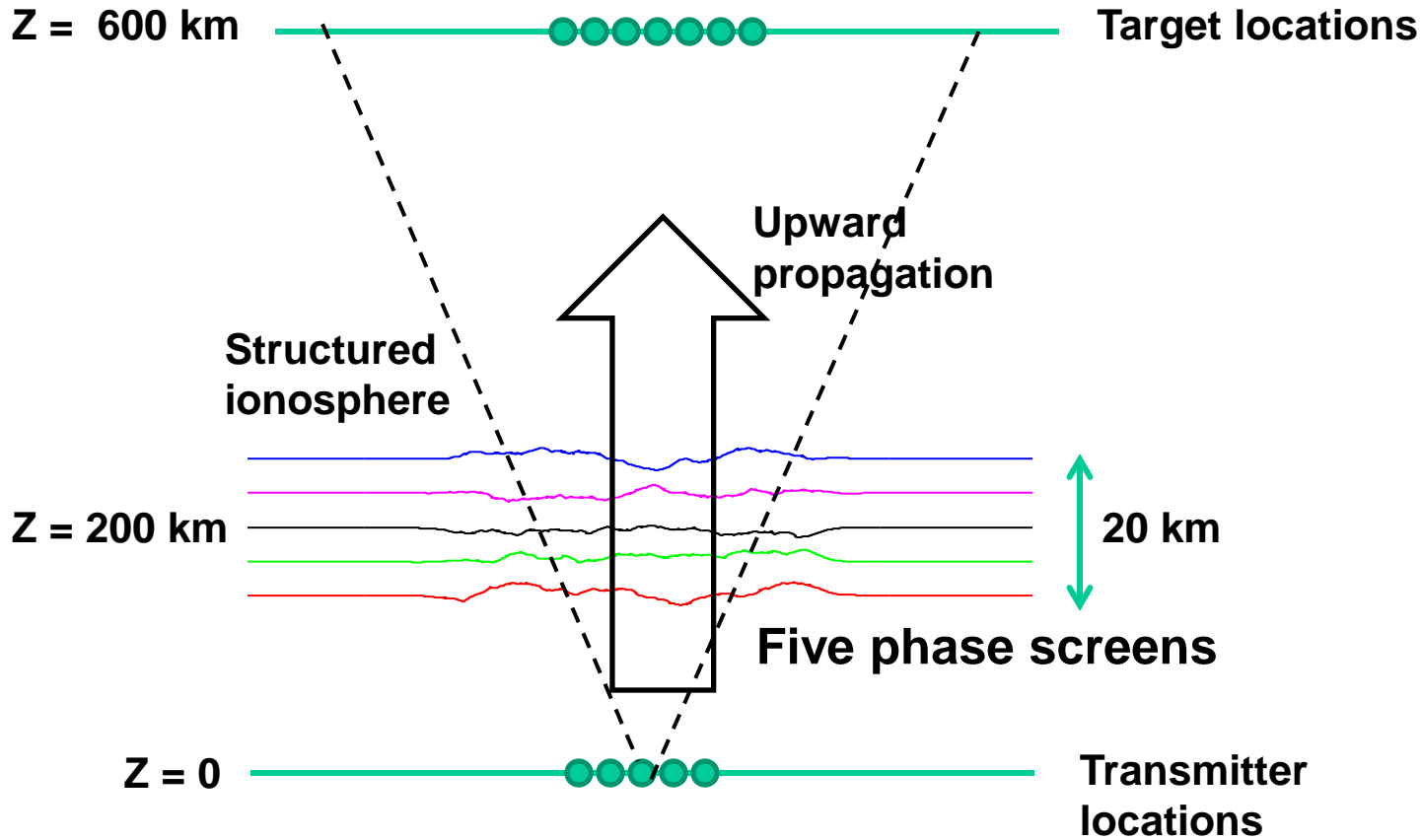
$$U(\theta, z'_2) = \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} U(\theta', z'_1) \exp\left\{\frac{i2\pi^2 q_\theta^2}{k} \left(\frac{1}{z'_1} - \frac{1}{z'_2}\right)\right\} \exp(i2\pi q_\theta(\theta - \theta')) dq'_\theta d\theta'$$

$$= \sqrt{\frac{iz^*}{\lambda}} \int_{-\infty}^{\infty} U(\theta', z'_1) \exp\left\{\frac{-ikz^*}{2}(\theta - \theta')^2\right\} d\theta'$$

**where**

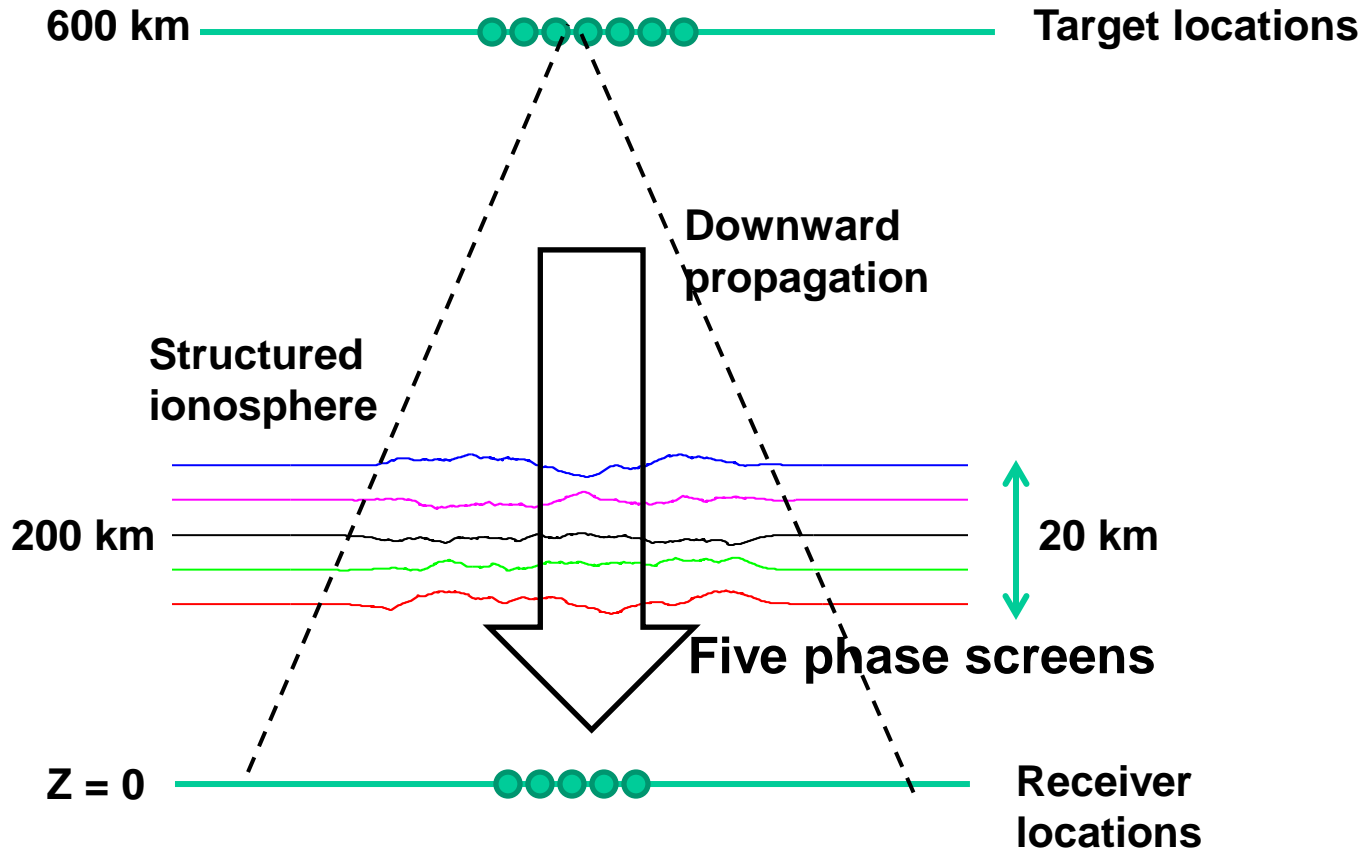
$$\frac{1}{z^*} = \left(\frac{1}{z'_1} - \frac{1}{z'_2}\right)$$

# Propagation Geometry Used in the Following Examples

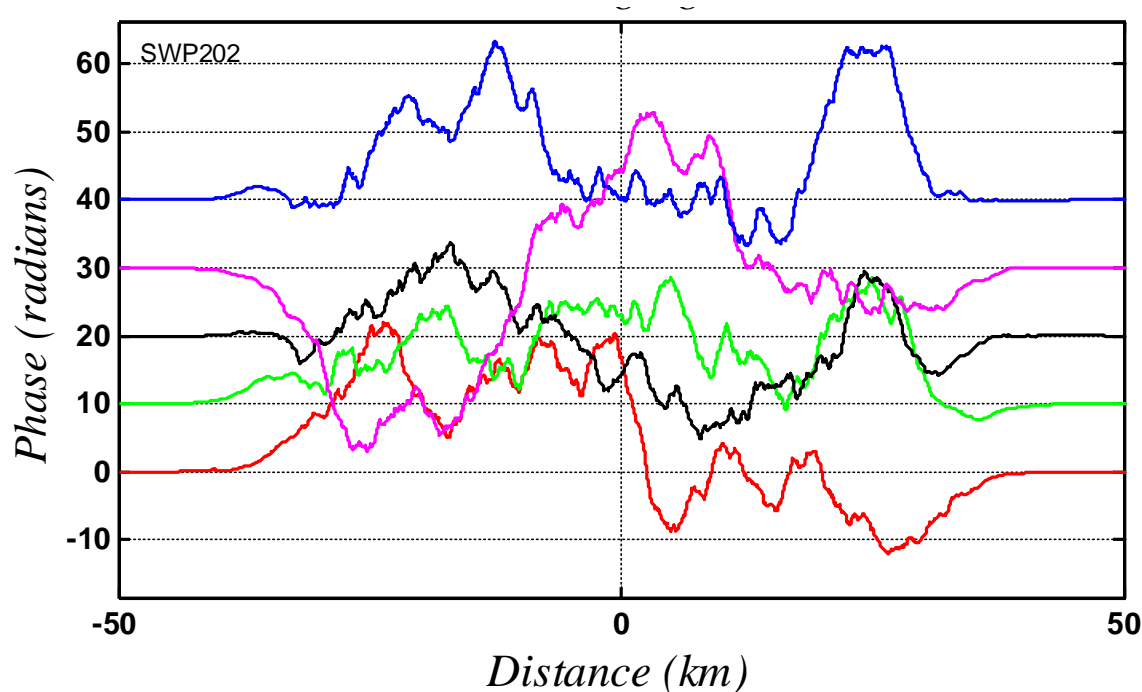




# Propagation Geometry Used in the Examples

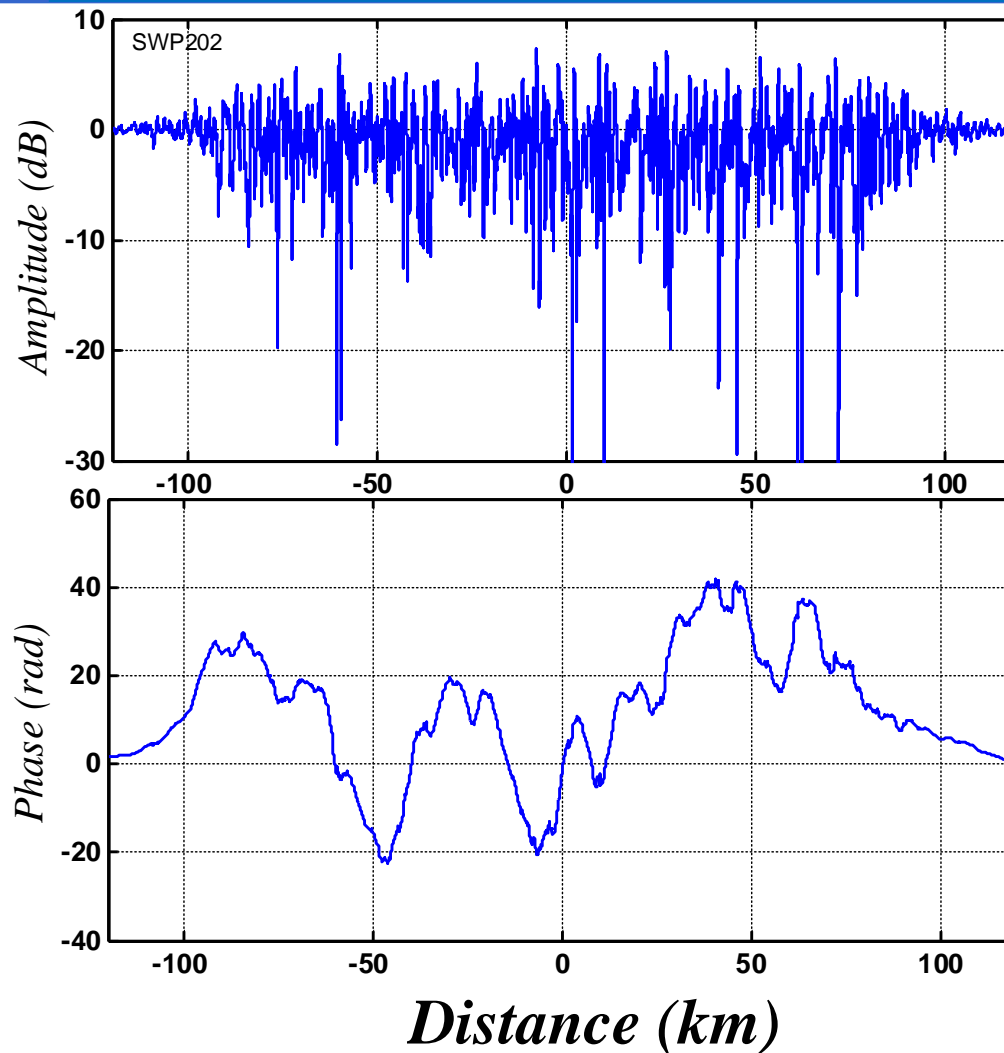


# Close-up of Five Phase Screens



- Values of phase shown are separated by 10 radians
- Screens extend in altitude from 190 to 210 km
- Length of phase screen at 200 km altitude is 200 km
- Phase screens are generated to have a  $K^{-3}$  PSD, outer scale of 5 km, inner scale of 10 m, and are comprised of  $2^{19}$  points.

# Electric Field in the Target Plane Due to a Single Transmitter



**Electric field at  $z = 600$  km caused by a single element located at  $z = 0$ , after propagation through five phase screens**

# Two-way Value of the Scintillation Index



**Definition of the S4 scintillation index, the normalized standard deviation of the received power**

$$S_4^2(\text{one-way}) = \frac{\langle (s - \langle s \rangle)^2 \rangle}{\langle s \rangle^2}$$

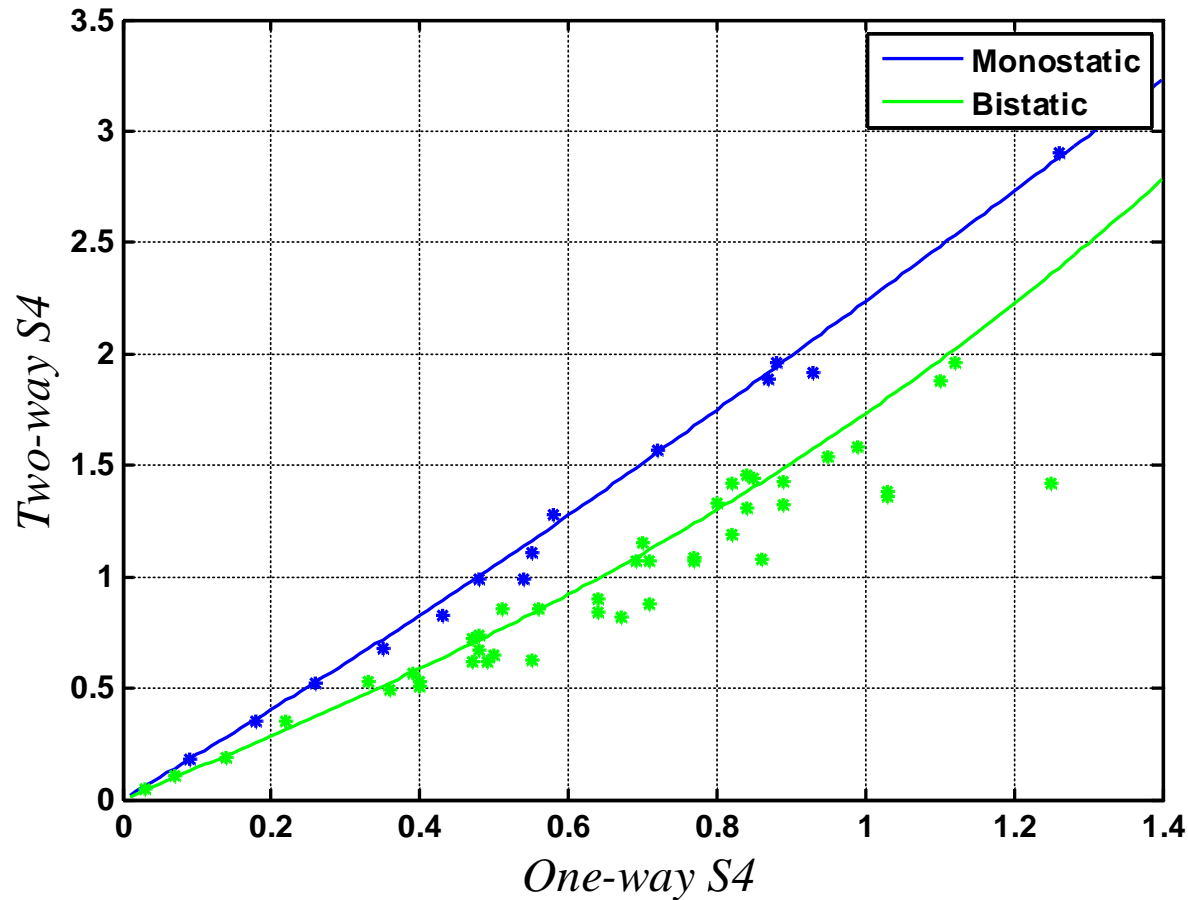
**For monostatic (radar) two-way propagation**

$$S_4^2(\text{monostatic}) = S_4^2(\text{one-way}) \left( \frac{4 + 6S_4^2(\text{one-way})}{1 + S_4^2(\text{one-way})} \right)$$

**For bistatic two-way propagation with independent up and down paths**

$$S_4^2(\text{bistatic}) = S_4^2(\text{up}) (2 + S_4^2(\text{up}))$$

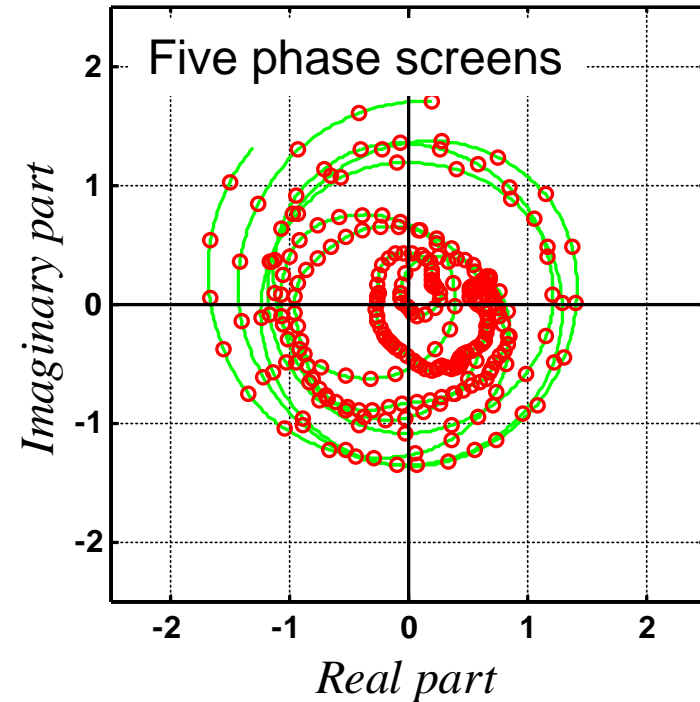
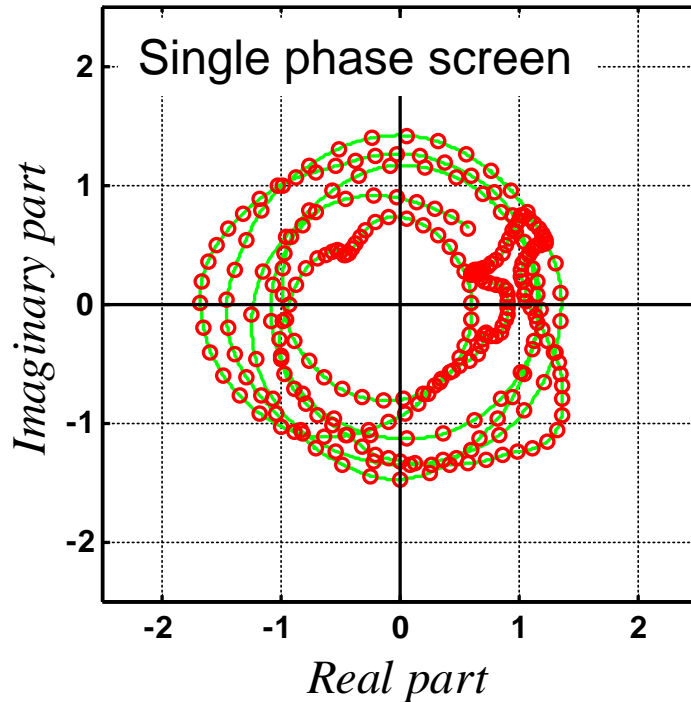
# Scintillation Index for Two-way Propagation



**Theory: solid lines; Simulation: dots**

**Radar detection performance is a strong function of  $S_4$**

# Reciprocity is Satisfied



- **Reciprocity: Field is same if transmitter and receiver are interchanged.**
- **The figures show I/Q plots of the complex one-way field comparing upward (green curve) and downward (red circles) propagation**
- **Upward propagation from single transmitter to many receive locations. Downward propagation from original receive locations to the single original transmitter location**



## Field at target plane due to many transmitter elements

$$V(x, z_{tar}) = \sum_{xmtr} V(x_{xmtr}, z_{xmtr}) G_{up}(xmtr \rightarrow tar)$$

## Field at receiver plane due to scatterers in target plane

$$\begin{aligned} V(x, z_{rcvr}) &= \sum_{tar} V(x, z_{tar}) G_{down}(tar \rightarrow rcvr) \\ &= \sum_{xmtr} V(x_{xmtr}, z_{xmtr}) \sum_{tar} G_{up}(xmtr \rightarrow tar) G_{down}(tar \rightarrow rcvr) \end{aligned}$$

## Following two examples of two-way propagation:

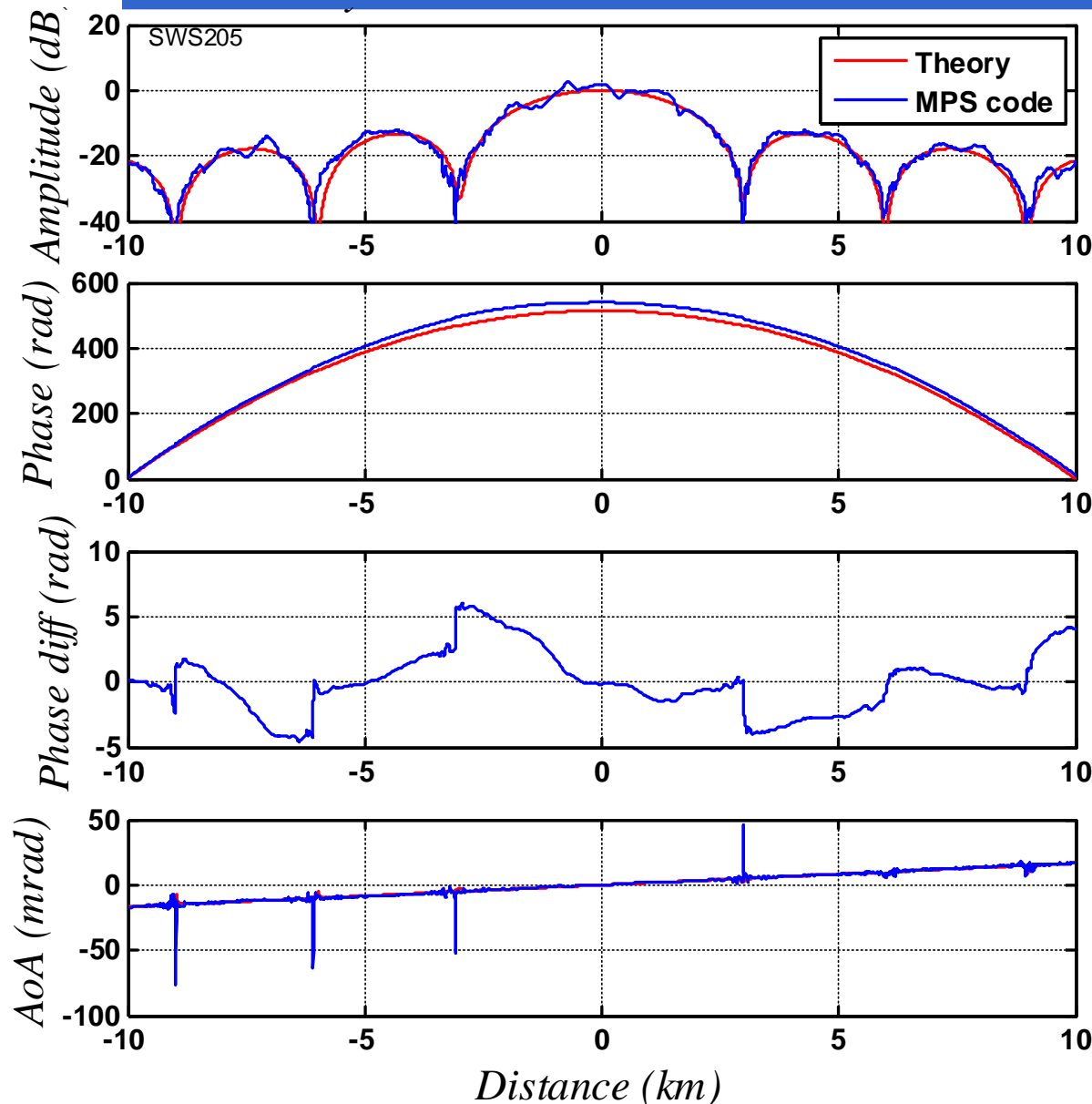
**One transmitter at center of MPS grid**

**Upward propagation through five phase screens**

**401 target scatterers at  $z = 600$  km, spaced by  $\lambda/2$**

**Downward propagation back to receiver plane**

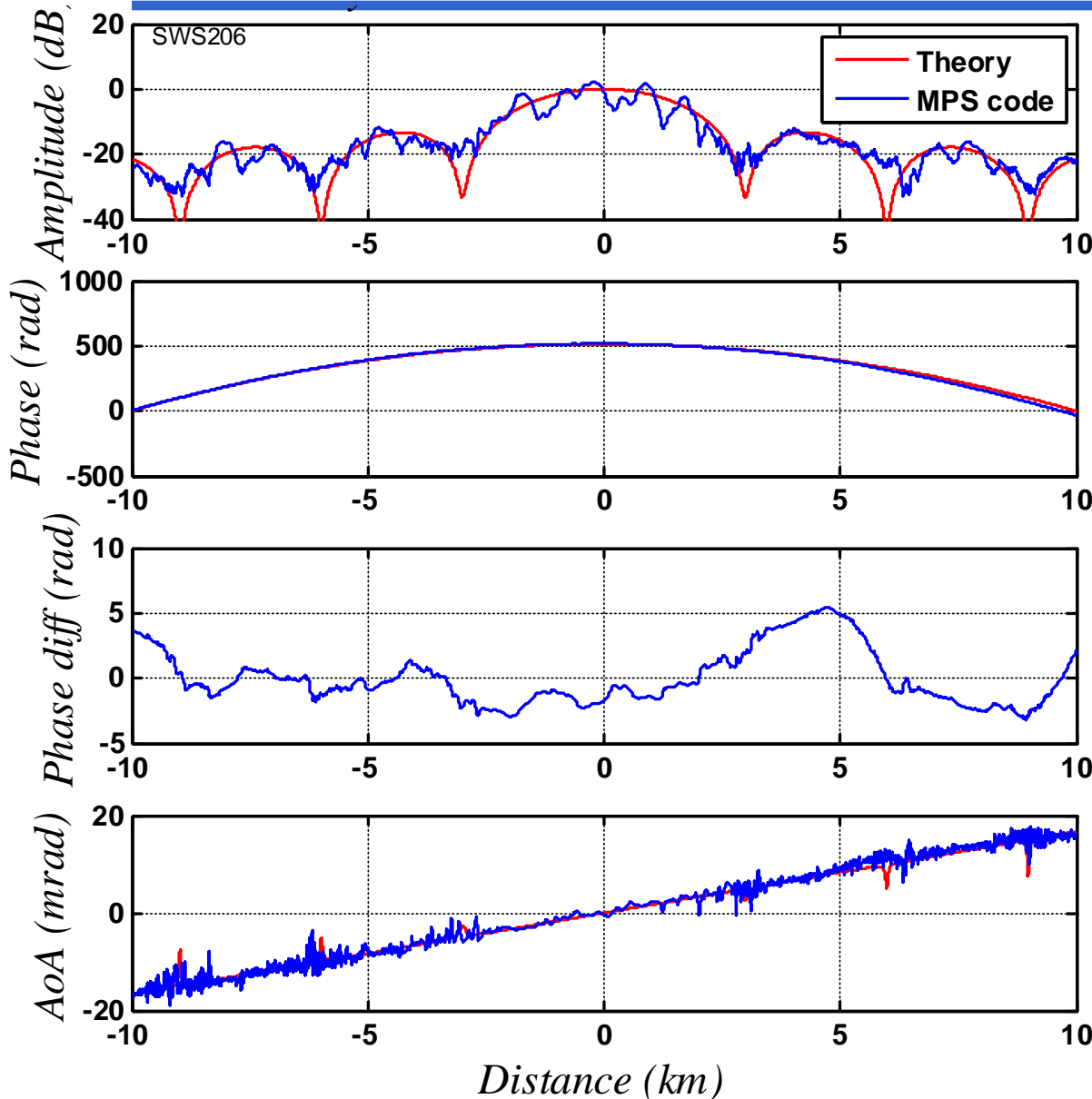
# Two-way Propagation, Weak Scattering, Linear Group of Scatterers



- 5 screens near  $z = 200$  km
- 401 scatterers at  $z = 600$  km
- $S_4(\text{one-way}) = 0.16$
- Figure shows small portion of MPS grid
- Smooth red curve is theory for case of no scintillation
- Blue is MPS result
- Measurement of AoA uses 10-m antenna & correlation technique



# Two-way Propagation, Stronger Scattering, Linear Group of Scatterers



- 5 screens near  $z = 200$  km
- 401 scatterers at  $z = 600$  km
- $S_4(\text{one-way}) = 0.46$
- Figure shows small portion of MPS grid
- Smooth red curve is theory for case of no scintillation
- Blue is MPS result
- Measurement of AoA uses 10-m antenna & correlation technique

# Conclusions

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- **Originally developed for application to synthetic aperture radar**
- **Includes the correlation of signals propagating on closely-spaced paths**
- **Avoids the small-scene approximation**
- **Code design allows for variation in RCS of the target scatterers**
- **Additional but straightforward work needed for:**
  - **3D propagation**
  - **Application to wide bandwidth waveforms**